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# Single Silicon Vacancy Centers in 10 nm Diamonds for Quantum Information Applications

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Supporting Information

ABSTRACT: Ultrasmall (about 10 nm), low-strain, artificially produced diamonds with an internal, active color center have substantial potential for quantum information processing and biomedical applications. Thus, it is of great importance to be able to artificially produce such diamonds. Here, we report on the high-pressure, high-temperature synthesis of such nanodiamonds about 10 nm in size and containing an optically active, single silicon-vacancy color center. Using special sample preparation technique, we were able to prepare samples containing single nanodiamonds on the surface. By correlating



atomic-force microscope images and confocal optical images, we verified the presence of optically active color centers in single nanocrystals, and using second-order correlation measurements, we proved the single-photon emission statistics of these nanodiamonds. These color centers have nonblinking, spectrally narrow emission with narrow distribution of spectral width and positions of zero-phonon line, thus proving the high quality of the nanodiamonds produced.

KEYWORDS: SiV centers, nanodiamonds, diamond growth, photons, single photon sources, high pressure high temperature, strain

# INTRODUCTION

Single-photon sources are of great interest for various applications in quantum information<sup>1</sup> and in particular quantum cryptography.<sup>2</sup> In addition, color centers in diamond attract a lot of attention due to the possibility of generating single photons at room temperature, low spread in spectral characteristics, and access to electron spin.<sup>3,4</sup> Beyond singlephoton applications, ultrasmall nanodiamonds containing color centers are of great interest for bioimaging<sup>5</sup> and biosensing.<sup>6,7</sup>

Among the most practical and interesting color centers in diamond are nitrogen-vacancy (NV) centers,<sup>8</sup> nickel-based (NE8) centers,<sup>9</sup> nickel-silicon complexes,<sup>10</sup> chromium-related color centers,<sup>11</sup> germanium-vacancy color centers,<sup>12</sup> and tinvacancy color centers,<sup>13</sup> as well as silicon-vacancy (SiV)

centers.<sup>14</sup> The most developed is NV center, which has properties that include an optical readout of spin state and a long coherence time up to 1.8 ms,<sup>15</sup> which are the focus of interest for this center.<sup>3</sup> NV centers are available in singlecrystal diamonds<sup>16</sup> and nanodiamonds<sup>17,18</sup> and can be produced easily during the growth process or by ion implantation.<sup>19</sup> Huge progress in quantum information processing has been achieved with NV centers: two-qubit quantum gates,<sup>20</sup> quantum register based on NV spins,<sup>21</sup> quantum memory at room temperature exceeding 1 s,<sup>22</sup> and

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Figure 1. (a) Structure of the SiV center with a silicon atom in a split vacancy site. (b) Schematic depiction of an assembly of toroid-type highpressure device used for synthesis of nanosize diamond materials on the basis of the hydrocarbon metal catalyst-free growth systems: (1) reaction zone, (2) graphite heater, (3) catlinite container and catlinite thermoinsulating disks, (4) tungsten carbide anvils, (5, 6) steel supporting rings.



Figure 2. (a) SEM image of the diamond material in the middle of the reaction zone right after growing. (b) Histogram of the size distribution of synthesized nanodiamonds obtained from TEM measurements. (c) TEM image of diamond slurry. (d) HRTEM image of the single nanodiamond. (e) Raman spectra of nanodiamonds right after growing. (f) Raman spectra of nanodiamonds after cleaning and centrifugation.

much more. The NV center still has some drawbacks that limit its application, however. First, about 3% of its emission is concentrated in the zero-phonon line  $(ZPL)^{23}$  while the overall spectra is 200 nm wide. In addition, the structure of the NV center is similar to the polar molecule, having high polarizability, making the NV center very sensitive to surrounding defects.<sup>24</sup> While this property gives rise to various sensor applications of NV center,<sup>7,25,26</sup> it considerably complicates any quantum information applications.

Recently, SiV centers were suggested as an alternative to NV centers<sup>27–30</sup> due to their strong ZPL containing about 70% of emission<sup>31</sup> and narrow width of around 5 nm at room temperature.<sup>14</sup> Another advantage is their emission at 738 nm in a spectral region where the background fluorescence of the surrounding material is weak. The first single-photon experiments on SiV centers were conducted with color centers created by ion implantation in IIa-type diamonds;<sup>14</sup> these,

however, revealed unfavorably low, single-photon emission rates on the order of only 1000 counts/s, despite their short lifetime of 1.2 ns. In the case of nanodiamonds or shallow implanted color centers, the luminescence of SiV centers may be quenched by the presence of strain<sup>32</sup> or other impurities,<sup>33</sup> thus underscoring the need for high-purity diamonds. Recently, implanted SiV in pure diamond films has shown reasonable count rates on the order of 200 kc/s, comparable with other color centers in diamond.<sup>34</sup> Similar results were obtained in as-grown SiV center in bulk plates.<sup>35</sup> Unfortunately, even in these bright SiV centers, the radiative quantum yield of SiV color centers at room temperature was still relatively low. For example, in recent experiments for SiV centers in CVD films, the radiative quantum yield was found to be only 0.05 due to the presence of decay via optical phonons,<sup>36</sup> which makes it difficult to use directly in quantum

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information application and requires strong enhancement overcoming low radiative yield.<sup>37</sup>

An important parameter that largely defines the application of the color center in diamond is the size of the crystal. Applications as biosensing and bioimaging require molecularsize resolution<sup>5</sup> that powerfully constrains nanodiamond size. Furthermore, quantum information applications can benefit a lot from ultrasmall diamonds with a single SiV center inside. Here, the small size helps prevent scattering of light on the nanodiamond containing the SiV center and suppresses the phonons that cause decoherence. Ultrasmall nanodiamonds containing SiV centers were successfully found in meteorite nanodiamonds of ultrasmall (below 5 nm) size.<sup>38-40</sup> Chemical vapor deposition demonstrated the possibility of creation of sub 10 nm nanodiamond, containing SiV color centers but exhibiting some agglomeration.<sup>41</sup> Just recently considerable advance was reached by combining chemical vapor deposition with oxygen etching.<sup>42</sup> The method enabled production of sub 10 nm diamonds, contacting around 3 color centers per nanocrystal with 6.4 nm ZPL. In this paper, we demonstrate the growing of ultrasmall (about 10 nm) nanodiamonds containing a single or a few bright SiV color centers with spectrally narrow ZPL line (down to 4.6 nm for single color center and on average 5.9 nm per crystal) using the highpressure, high-temperature method (HPHT).

#### EXPERIMENTAL SECTION

Nanodiamond Synthesis. Studied nanodiamonds with SiV color centers were synthesized by a metal-free HPHT method on the basis of a hydrocarbon-containing growth system.<sup>43–45</sup> Naphthalene  $C_{10}H_8$ (Chemapol) and tetrakis(trimethylsilyl)silane  $C_{12}H_{36}Si_5$  (Stream Chemicals) were used as the main hydrocarbon and silicon-doping components of the growth mixtures. Cold-pressed pellets of the initial homogeneous mixture (5 mm in diameter and 3 mm in height) were inserted into a graphite container that also served as a heater. HPHT treatment of the samples was carried out in a high-pressure, "toroid"type apparatus<sup>46</sup> (see Figure 1b). The experimental procedure consisted of loading the high-pressure apparatus to 8.0 GPa and heating the samples up to temperature of diamond synthesis (1300-1400 °C) with short (about 1s) isothermal exposure at this temperature. X-ray diffraction (XRD) and Raman spectroscopy (see Supporting Information for details) and transmission (TEM) electron microscopies were used for preliminary characterization of synthesized diamond materials. According to the obtained data, the size fraction of the diamond (ultrasmall, about 10 nm, nanometer up to 100 nm, submicrometer 100 to 1000 nm, and micrometer, more than 1000 nm) in the conversion products is determined basically by the composition of initial growth mixture, temperature, and synthesis time.<sup>44</sup> Adding fluorocarbon compounds (fluorinated graphite  $CF_{1,1}$ ) to the initial growth mixtures decrease the threshold of the diamond formation temperature.<sup>43</sup> Therefore, synthesis at the closest point to the temperature threshold of the beginning of the formation of diamond and short (about 1 s) time of isothermal exposure led to the increase of the ultrasmall fraction of diamond material in the conversion products, which can reach 50% of the total yield. In this case, taking into account the radial temperature gradient in the reaction zone of the high-pressure apparatus with the use of a cylindrical electrical resistivity graphite heater, the maximum content of an ultrasmall sized diamond fraction is observed in the middle of the reaction zone (see Figure 2a). The maximum temperature reaches the zone of contact of the material with the walls of a graphite heater; therefore this area is characterized by a high content of monocrystals with submicrometer and micrometer-size fractions.

**Sample Preparation.** Ultrasmall nanodiamonds have a strong tendency toward aggregation, and in order to prevent clustering, we performed the following procedure. First, the glass coverslip (Menzel Gläser) with a preexisting gold mask was cleaned in a piranha solution

(H<sub>2</sub>SO<sub>4</sub> and H<sub>2</sub>O<sub>2</sub> in a ratio of 3 to 1, respectively) and kept at 120 °C for 1 h. Then, the slurry with nanodiamonds was sonicated for 1 h at a power of 170 W. For sonication the Ultrasonic Cleaner CD-4820 was used. Finally, 10  $\mu$ L of slurry was dropped onto the coverslip and evaporated in the vacuum chamber using a Harrick Plasma PDC-002 plasma cleaner.<sup>47</sup>

**Confocal Setup.** To analyze the optical properties of SiV centers, we used a homebuilt confocal setup (Figure 3a) with a pumping diode



**Figure 3.** (a) Schematic of the confocal microscope. (b, c) Lifetime measurement for SiV center. Fast decay during the first nanoseconds is due to background fluorescence.

laser at 532 nm (Coherent Compass 300) for exciting the SiV center in nanodiamonds, and galvo mirrors (Cambridge Technologies) for scanning the specimen. The emission was collected with a high numerical aperture oil objective (NIKON Apo TIRF 100X NA 1.49) with a working distance of 120  $\mu$ m. The assembled fluorescence passes through the combination of the long-pass optical filter with cutoff at 600 nm and notch filter with a stop-band centered at 532 nm to reject the excitation light and Raman signal from the collected emission. The luminescence that passed through the filters was coupled into a single-mode fiber (Thorlabs SM600) and guided to a Hunbury-Brown-Twiss (HBT) interferometer that consisted of two avalanche photodiodes (PerkinElmer SPCM-AQRH-14-FC) and a 45:55 beam splitter or directed to the homebuilt spectrometer with monochromator (Jarrell Ash model 82-422) and CCD camera (Starlight Xpress SXV-H9C). A time-correlated, single-photon counting module (Picoquant Picoharp 300) was used to obtain second-order photon correlation histograms from SiV centers.

For the lifetime measurements, a picosecond diode laser (Picoquant LDH-P-FA-530 XL) was used with a same timecorrelation, single-photon counting module. Due to the fact that the lifetime of the SiV center excited state is longer than the length of the laser pulse, the fluorescence signal recorded after the applied laser pulse should exhibit exponential decay. In fact, the actual signal observed consisted of two exponents: fast and slow (Figure 3b and Figure 3c). If the fast exponential component corresponds to background fluorescence signatures that do not have a slowly decaying exponential component, then the slow decay "tail" is associated with the lifetime of SiV center in an excited state.

With the exception of the spectrometer measurements, the emission was additionally filtered by a  $737 \pm 10$  nm narrowband filter.

## RESULTS AND DISCUSSION

The SiV center consists of a silicon atom and a lattice vacancy. The silicon atom, which substitutes a carbon atom, relaxes its lattice position toward two neighboring vacancies. This composition existed along  $[111]^{48}$  crystallographic axis in



Figure 4. (a) Confocal map of the sample with deposited nanodiamonds. Red circles indicate nanodiamonds with optical active SiV centers. (b) AFM image of the area, marked with red square in confocal image. The blue line demonstrates direction along which 1D image was taken to determine particle size. (c) Size distribution of the nanodiamonds containing optical active SiV centers. (d) 1D profile of nanodiamond aiding in the visualization of nanodiamond size.

diamond (see Figure 1a). The defect is negatively charged and belongs to the  $D_{3d}$  group of symmetry.

The scanning electron microscope (SEM) image of the diamonds formed in the central part of the reaction zone is presented in Figure 2a. One can see a large amount of nanodiamonds (with sizes well below 100 nm) with a considerable fraction of ultrasmall (<10 nm and about 10 nm) diamonds and a small quantity nondiamond, carbon materials.

Ultrasmall fraction was separated using the following procedure. First, the samples were treated in a 40% solution of hydrogen peroxide for chemical purification of the diamond material. Then, separation of the ultrasmall fraction of diamond was carried out in several stages that consisted of ultrasonic dispersing of diamond particles in an aqueous medium in a UP200Ht dispersant (Hielscher Ultrasonic Technology) for 1 h at a power of 200 W and subsequent three-stage centrifugation of aqueous dispersions of diamond materials at 1000g, 3000g, and 5000g for 5 min each time. As a result, the diamond fraction with the dominant content of particles of about 10 nm was obtained. A transmission electron microscope (TEM) image of this diamond fraction (see Figure 2c), high resolution TEM image of the single nanodiamond (see Figure 2d), and histogram of the size distribution calculated from this are presented in Figure 2b (total of 149 nanodiamonds).

We performed the Raman measurements for nanodiamonds right after growing (see Figure 2e) and after cleaning and centrifugation procedure (see Figure 2f). The reduction of the Raman line in range  $1500-1800 \text{ cm}^{-1}$  is observed. This reduction indicates reduction of non-diamond carbon containing materials, which stayed during growth. After cleaning and centrifugation we observe small peak at 1590 cm<sup>-1</sup> which was related to the G-band and caused by in plane vibrations of graphitic carbon.<sup>49</sup> Therefore, we can conclude about the dominating of the sp<sup>3</sup> diamond Raman line near 1332 cm<sup>-1</sup> in the grown nanodiamonds (see Supporting Information for more details). The ultrasmall nanodiamonds fraction was retained in the aqueous dispersion after centrifugation.

One of the interesting features of ultrasmall nanodiamonds is their ability to form aggregates. The use of water as a basis for nanodiamond slurry is known to lead to aggregation during the seeding procedure.<sup>50</sup> To avoid this, we used a vacuum evaporation seeding method<sup>47</sup> (detailed in Experimental Section), which allowed us to prevent aggregation during evaporation and observe SiV centers in single nanodiamonds of ultrasmall sizes.

To get insight into the size of the seeded aggregates containing optically active color centers, we used preprinted gold features on the glass coverslips under investigation (see Figure 4). This allowed us to correlate images taken with a confocal microscope (see Figure 4a) and with an atomic force microscope (AFM) (see Figure 4b; see ref 51 for more details). The AFM images were taken using an NTEGRA Spectra-M setup with silicon cantilever (NT-MDT HA NC) operated in taping mode. In total, 3 AFM maps of  $10 \times 10 \ \mu m^2$  $(256 \times 256 \text{ points scan with lateral resolution about 40 nm})$ and resolution in the normal direction to the scanned surface was 0.1 nm) were correlated with confocal maps. The measured sizes of the nanodiamond having color centers are presented in Figure 4c,d. From this measurement, it was verified that the distribution of nanodiamonds containing the SiV color centers (see Figure 4c) has median  $9.3 \pm 2.5$  nm, which is very similar to the distribution of the nanodiamond sizes themselves (see Figure 2b), whose median is  $7.5 \pm 5$  nm. Besides, the probability of finding at least one SiV center in single nanodiamond is about 0.5%, whereas the probability of finding single SiV center in single nanodiamond is approximately 0.15%. Thereby, the probability of finding nano-



**Figure 5.** (a) Second-order autocorrelation function obtained from a single SiV center. The inset illustrates measurement of detector's jitter and its approximation by Gaussian distribution. The green horizontal lines show  $g_2(0)$  for 2, 3, and 4 equally bright emitters, assuming zero noise level. (b) Saturation curves for a single SiV centers. The red, orange, and green solid lines represent the fit to the saturation; the dashed blue line is the background contribution to the counts detected. (c) Time trace of SiV center under various excitation powers. (d) Spectra obtained for a single SiV center, full width half maximum = 4.6 nm.

diamond with single SiV is approximately 30% among nanodiamonds with SiV color centers (see Supporting Information for more details). Thus, color centers were observed in single nanodiamonds or just few nanodiamonds per aggregate.

The optical properties of the nanodiamonds were studied using a homebuilt, confocal microscope with an immersion oil objective of NA 1.49. For the excitation source, we used a 532 nm diode laser (see Experimental Section for details). The relative brightness of the nanocrystals most likely is related to the number of color centers per nanocrystal. To prove that we could grow nanocrystals containing a single color center with this method, we performed second-order correlation function  $g_2(\tau)$  measurements (see Figure 5a). The depth of  $g_2(0)$ characterizes the number of emitters detected, while the width of  $g_2(\tau)$  is limited by the color center's excited state lifetime and can only decrease as excitation power increases.<sup>52</sup> For nanocrystals, we measured the lifetime of 1.41 ns (see Figure 3b and Figure 3c), while the time resolution of a pair of our detectors was measured as 672 ps (see inset in Figure 5a). Therefore, the time resolution of the detecting system considerably limits the possible depth of  $g_2(\tau)$  at 0 time delay  $\tau$ .<sup>52</sup> To correct for this distortion, we performed a deconvolution procedure (see Figure 5a and Supporting Information for more details). After deconvolution, we obtained the achieved  $g_2(0)$  value, 0.41, which was below the threshold of 0.5 for two equally bright color centers.

Color center brightness depends strongly on the efficiency of excitation reaching maximum value at saturation of the center. The absorption maximum for a SiV color center is at ZPL, making it relatively hard to reach experimentally complete saturation of the center emission with green excitation.<sup>31</sup> Nevertheless, information on saturation counts could be extracted from analysis of the saturation curve (Figure 5b). To analyze the saturation curve, we used the standard model for fitting:

$$I = \frac{I_0 p}{p_0 + p} + \alpha p \tag{1}$$

where  $I_0$  is the counts in saturation, p is the laser power,  $p_0$  is the laser power in saturation, and  $\alpha$  is the coefficient for a linear substrate noise. Here, we analyzed the emission of 7 single color centers in our diamonds. By fitting these experimental data, we found that a single SiV center in a nanocrystal has an average saturation count rate of 75  $\pm$  47 kcounts/s at the saturation power  $3.2 \pm 2.4$  mW; these values were gained from fit's extrapolation because reaching the saturation with green laser can be relatively hard. The whole error is square root from statistical measurements and fit error. The observed saturation behavior greatly varied from center to center (Figure 5b), with the fitted  $I_0$  values ranging from 19  $\pm$ 6 kcounts/s to 143  $\pm$  30 kcounts/s and with the  $p_0$  ranging from  $0.6 \pm 0.3$  mW to  $6.9 \pm 2.3$  mW. This intensity level is comparable with SiV centers in CVD nanodiamonds<sup>53</sup> and HPHT nanodiamonds,<sup>54</sup> as grown in bulk diamond plate<sup>35</sup> and ion implantation followed by annealing.<sup>34</sup> Remarkably, this high count rate is accompanied by rather stable emission, depicted in Figure 5c.

In addition, we performed detailed spectral studies for various nanodiamonds ranging from nanocrystals containing a single SiV to nanocrystals or aggregates containing ensembles with approximately 10 SiV centers. The approximate number of color centers was estimated as a ratio of saturation count to the single color center average number of counts in saturation.

Figure 5d shows the typical spectrum of SiV centers in synthesized nanodiamonds. We measured a total of 24 spectra for SiV centers, contained in nanocrystal found from correlated confocal and AFM maps. Besides, spectra of extra 11 nanocrystals were recorded, for which correlation measurements were not performed. Each spectrum was fitted with the sum of 3 Lorentzian functions. One of the Lorenzians approximated ZPL and the rest, phonon sideband. We found



Figure 6. (a) Histogram of ZPL line position. (b) Histogram of fwhm of ZPL lines of SiV centers observed from 1-10 SiV centers per crystal.

that the ZPL position has a narrow distribution around 738.06 nm, with a standard deviation of 0.27 nm (Figure 6a). This narrow distribution indicates the low level of strain in our nanodiamond, because ZPL line position in general depends on the strain.55 We note that the spread of the ZPL line position for SiV centers grown by other methods is considerably higher, as previously indicated: in the case of CVD nanodiamond synthesis,<sup>56</sup> it is about 1.5 nm. For HPHT nanodiamonds grown with other conditions ZPL line position is 737  $\pm$  10 nm.<sup>54</sup> For ultrasmall nanodiamonds found in natural meteorite, ZPL position was found to be 735.7 with variations of  $\pm 5$  nm.<sup>38</sup> At the same time, the width of ZPL line at room temperature has a mean of 5.9 nm with a standard deviation of 0.8 nm (Figure 6b). This is a rather standard value determined mostly by the strong coupling of excited levels of SiV color centers via the phonons in diamond.<sup>57</sup> This mixing is obviously temperature dependent<sup>58</sup> and thus should only be compared at a known temperature. Using the Raman shift measurements for diamond material right after growing and the ZPL line position statistics, we assessed the strain level in produced nanodiamonds to be less than 0.7 GPa (see Supporting Information for details).

## CONCLUSIONS

We synthesized ultrasmall nanodiamonds (smaller than 10 nm) containing single or few SiV color centers. The nanodiamonds demonstrate high single-photon brightness (>100 000 cps) at room temperature. Both the nanocrystals containing a single color center and those with relatively large concentrations of SiV centers (up to 10 SiV per 10 nm crystal) have narrowband luminescence and extremely stable ZPL position, which indicate strain level less than 0.7 GPa and correspondingly low number of defects, different from SiV centers s in the nanodiamonds grown. This high quality, ultrasmall nanodiamond may be of great interest for quantum information and bioimaging applications.

## ASSOCIATED CONTENT

#### **S** Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsanm.9b00580.

Calculation of probability of finding SiV in single nanodiamonds, size measurements from AFM maps, Raman measurements and preliminary characterization of grown diamond material, strain assessment by the ZPL position analysis, measurement of the setup jitter function, and deconvolution procedure (PDF)

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## **Author Contributions**

The main contributors of this paper are Stepan V. Bolshedvorskii and Anton I. Zeleneev, who performed the majority of the optical measurements and developed and realized correlation of confocal and AFM maps, and Valery A. Davydov who was responsible for diamond synthesis. Alexey V. Akimov developed the initial idea of this work and supervised all parts of the study. Vadim V. Vorobyov, Vladimir V. Soshenko, and Olga R. Rubinas helped with acquiring and analysis of optical data. Leonid A. Zhulikov prepared the samples and analyzed nanodiamonds positioning on the sample. Pavel A. Pivovarov performed AFM measurements. Vadim N. Sorokin and Andrey N. Smolyaninov supervised optical and AFM investigation process. Valery A. Davydov and Liudmila F. Kulikova performed full sample diamond material synthesis and purification procedure. Sergey G. Lyapin performed Raman measurements of full sample diamond material. Anastasiia S. Garanina, Viatcheslav N. Agafonov, and Rustem E. Uzbekov performed separation of the ultrasmallsized fraction of diamonds and transmission electron microscope investigation of the diamond materials for optical experiments.

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#### Notes

The authors declare no competing financial interest.

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